



# Fast determination of phosphorus concentration in phosphogypsum waste using calibration-free LIBS in air and helium

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## Abstract

This work represents a novel method to determine phosphorus (P) concentration in phosphogypsum (PG) waste samples using calibration-free laser-induced breakdown spectroscopy (LIBS). A 50 mJ Q-switched Nd: YAG laser has generated the PG LIBS spectrum. Spectroscopic analysis of plasma evolution has been characterized by electron density  $N_e$  and electron temperature  $T_e$  using the emission intensity and Stark broadening for P I characteristic lines 213.61, 214.91, and 215.40 nm under non-purged (air) and purged (helium) conditions. It was found that both  $T_e$  and  $N_e$  have significant changes linearly with P concentrations 4195, 5288, 6293, and 6905 ppm. The plasma  $T_e$  and  $N_e$  values increased from about 6900 to 10,000 K and  $1.1 \times 10^{17}$  to  $3.4 \times 10^{17} \text{ cm}^{-3}$ , respectively, for the non-purged PG. On the other hand,  $T_e$  and  $N_e$  ranged from 8200 to 11,000 K and  $1.4 \times 10^{17}$  to  $3.5 \times 10^{17} \text{ cm}^{-3}$ , respectively, for the PG purged with helium. It is concluded that  $T_e$  and  $N_e$  values represent a fingerprint plasma characterization for a given P concentration in PG samples, which can be used to identify P concentration without a PG's complete analysis. These results demonstrate a new achievement in the field of spectrochemical analysis of environmental applications.

**Keywords** Phosphorus · Phosphogypsum · Laser-induced breakdown spectroscopy · Helium · Calibration-free · Electron temperature · Electron density

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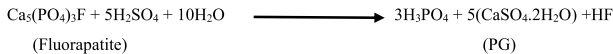
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## 1 Introduction

One of the significant nutrients is phosphorus (P), as well as nitrogen (N) and potassium (K), which are essential for plant growth and crop production. Phosphorus deficiency in agricultural soil can be overcome using phosphate bearing fertilizers (Kaski et al. 2003). Phosphate ore is a crucial source for phosphate fertilizers and phosphorus for agricultural applications. Phosphate fertilizers' life cycle starts with phosphate ore (mainly of sedimentary or igneous origin) that is chemically treated with sulphuric acid to produce phosphoric acid and phosphogypsum (PG) ( $\text{CaSO}_4$ ) as a by-product waste. PG waste has considerable concentrations of radioactive materials ( $^{238}\text{U}$ ,  $^{226}\text{Ra}$ ,  $^{210}\text{Pb}$ ,  $^{232}\text{Th}$ , and others), phosphorus pentoxide ( $\text{P}_2\text{O}_5$ ), and some trace elements such as chromium, fluoride, zinc, antimony, arsenic, lead, cadmium, and copper (Calin et al. 2015; Cuadri et al. 2014; Potiriadis et al. 2011). So, a precise determination of phosphate content in fertilizers and PG is needed for agricultural good-practice and environmental protection.

Phosphogypsum (PG) is a by-product waste of phosphoric acid ( $\text{H}_3\text{PO}_4$ )/phosphate fertilizers manufacture where phosphate is attacked by sulfuric acid ( $\text{H}_2\text{SO}_4$ ) as shown in the following reaction equation (Calin et al. 2015; Cuadri et al. 2014; Potiriadis et al. 2011);



Although there are various valuable applications (such as agricultural and building materials), only about 15% of some billion tons of PG produced worldwide are recycled due to its impurities and contaminants (Nizevičienė et al. 2018; Pacheco-Torgal et al. 2012).

Recently Elisha J. J. investigated the elemental composition of nine randomly selected inorganic fertilizers using X-ray fluorescence spectrometric technique (XRF) (J. J. 2014). They analyzed some fertilizers purchased from Samaru Zaria for some toxic elements (Cr, Ni, Si, V, and Zn). Even though the obtained results confirm that the elemental concentrations were within the agricultural soils ranges, they failed to detect phosphate concentrations fluctuation. Raven and Loeppert evaluated trace elements and heavy metals in a wide variety of fertilizers and soil amendments by using both Atomic Absorption Spectrometry (AAS) and Inductively-Coupled Plasma Induced Emission Spectrometry (ICP-MS) (Raven and Loeppert 1997). İlknur ŞEN used near-infrared reflectance technique to analyse soil-fertilizer mixtures on-line (Şen 2003). The method was useful for use by farmers in the field without the need to collect samples and send them to the laboratory.

The use of high-power lasers as a spectroscopic analytical tool depends on vaporization, excitation, or ionized atoms/molecules from the material surface when a laser pulse is focused on it. The atomic emission spectroscopic study of emission lines from the induced plasma can give information about the investigated material's chemical composition. This technique of spectroscopic studies is known as Laser-Induced Breakdown Spectroscopy (LIBS). LIBS is typical atomic emission spectroscopy where a high power laser pulse strikes the sample surface to produce a plasma, which will then generate a unique fingerprint for the elemental content in the sample using their characteristics lines (Ahmed et al. 2020; Farooq et al. 2020; Farooq et al. 2018; Tawfik et al. 2015a).

LIBS is a promising technology for determining the elements in any material. It is a flexible approach that may be used in any environment. It can even be used in hostile environments (Cremers and Radziemski 2006). LIBS represents an advanced atomic analytical technique that overcomes the most disadvantages of the other traditional analytical techniques

like; atomic absorption spectrometry (AAS), atomic emission spectrometry (AES), and x-ray fluorescence spectrometry (XRF). Nevertheless, such methods are destructive techniques and necessitate time-consuming sample preparation procedures, while LIBS is non-destructive and needs no sample preparation (Qindeel and Tawfik 2014). So that, the samples used in LIBS are free from any progressive digestion, such as extraction chemicals or strong acids (Arantes De Carvalho et al. 2018; Fikry et al. 2021; Tawfik et al. 2015a; Wilberforce 2016). Moreover, LIBS is suitable for quick, real-time analysis (on-line), low cost, sensitive, and high spatial resolution analysis.

The sensitivity of LIBS technique is a critical factor, representing the enhancement of the detection limit of a minimal concentration of trace elements in the samples (Fikry et al. 2020b, 2020a; Rai et al. 2001; Yueh et al. 2002). Calibration free-LIBS analysis was proposed by Ciucci et al. (Ciucci et al. 1999), which can be applied to any element in the periodic table provided that the physics of the plasma process and local thermal equilibrium for the observed relevant spectral information under thin plasma conditions are experimentally validated (Tawfik and Mohamed 2007).

One of the most comfortable and cheapest techniques that has been utilized to enhance the sensitivity of LIBS system is purging the sample with inert gases. This improves the plasma electron temperature and density, which confirms the sensitivity and intensity of the induced breakdown spectrum (Kaski et al. 2003; Li et al. 2018).

Fertilizers from Saudi Arabia were analyzed using the LIBS technique (Farooq et al. 2012) to detect phosphorous, manganese, magnesium, molybdenum, iron, titanium, nickel, vanadium, calcium, cobalt, cadmium, tin, sulphur, aluminium, chromium, lead, and uranium. A relatively high concentration of health-hazardous elements such as Cd, Ni, and Pb were noticed. Gallou et al. (Gallou et al. 2009) used a portable LIBS system to detect elements such as F, Cl, P and U on-site. Using the LIBS technique, a quantitative method dedicated to determining phosphorus in fertilizers of different matrix compositions was developed. Phosphorus quantification in Brazil's organic and inorganic fertilizer samples was achieved with approximately 15% precision compared to other measurement techniques such as ICP (Marangoni et al. 2016). The LIBS technique was also used to assess agricultural products and precision agriculture compared to a standardized chemical analysis method. In their experiment, soil samples classified as thermic, fine-silty, and mixed from grassland and uplands in Tennessee (USA) were collected and smashed. The samples were analyzed using LIBS in the spectral range from 200 to 600 nm using a 532 nm Nd: YAG laser, 25 mJ pulse energy, 5 ns pulse width, with 10 Hz repetition rate. They stressed the necessity to enhance LIBS's measurement and prediction accuracy by improving the pre-treatment process, standard reference soil samples, and measurement method for a reliable quantification method (Cremers et al. 2001).

In this work, the calibration-free LIBS technique has been used to measure phosphorus concentrations in different phosphogypsum samples directly. In doing so, spectroscopic analysis of plasma evolution of laser-produced plasma has been characterized by elemental content spectra, plasma electron density, and electron temperature assuming the Local Thermodynamic Equilibrium (LTE) and optically thin plasma conditions under air (non-purge) and helium gas (purge) to enhance the LIBS signal.

## 2 Methodology

### 2.1 Sample preparation

The phosphogypsum samples were collected from the Rad Al-Khair chemical complex on the Arabian Gulf coast—Saudi Arabia. The samples were dried in an oven at 80 °C for about 48 h, then mechanically crushed in a ball mill, and sieved through a 1 mm mesh sieve. The samples were then pressed to produce tight, compact pellets with a diameter of 40 mm and a thickness of 10 mm under 200 MPa pressure for 5 min by a hydraulic oil press. The samples were pressed as pellets to decrease the sample's surface roughness, which enhanced the mass ablation's uniformity (Edwards and Winefordner 2007).

### 2.2 Analytical methods

Phosphogypsum samples were analyzed for various major and minor oxides' decomposition and their concentrations in the "Australian Laboratory Services' ALS Geochemistry limited company" labs. The analysis was done using inductively coupled plasma-atomic emission spectroscopy ICP-AES (ALS Geochemistry method ME-ICP06). In addition to that, the Thermal decomposition Furnace or TGA (OA-GRA05 or ME-GRA05) was performed.

### 2.3 LIBS setup

The laser-induced plasma (LIBS) system used in this study is the LIBS lab system (LTB Laser Technik Berlin GmbH) having the technical details shown previously elsewhere (Harmon et al. 2019; Ji et al. 2017; Merk et al. 2015). Briefly, in the current system, a Q-switched Nd:YAG laser (Quantel Ultra 50) beam with a 50 mJ laser pulses with pulse duration 5 ns at its fundamental wavelength 1064 nm, and adjustable repetition rate up to 30 Hz. The laser pulse energy was adjusted by changing the Q-switch delay using Sophi 64 bit software designed to control the spectrometer and the whole LIBS system. The light emitted by the plasma emissions at the focal volume was collected by a set of optics and focused on an optical fiber head. This fiber bundle delivered the light to an ARYELLE 200 butterfly spectrometer. ARYELLE 200 is an Echelle spectrometer with a pre-monochromator and active wavelength stabilization. The ARYELLE 200 provides a spectral resolving power 75,000–150,000 and spectral resolution from 13.8 to 37.5 pm over a wavelength range 192 – 750 nm displayable in a single spectrum in two spectral ranges (UV and Vis). The dispersed light from ARYELLE 200 was then detected by a gateable intensified charge-coupled detector ICCD (Andor I-Stare).

The emission spectra display, processing, and analysis were done using 2D and 3D Gram/32 software programs (National Instruments, USA). In addition to the atomic database used by the mentioned software's atomic database, spectral line identification was checked by the most up-to-date electronically published database (Kramida and Ralchenko 2018).

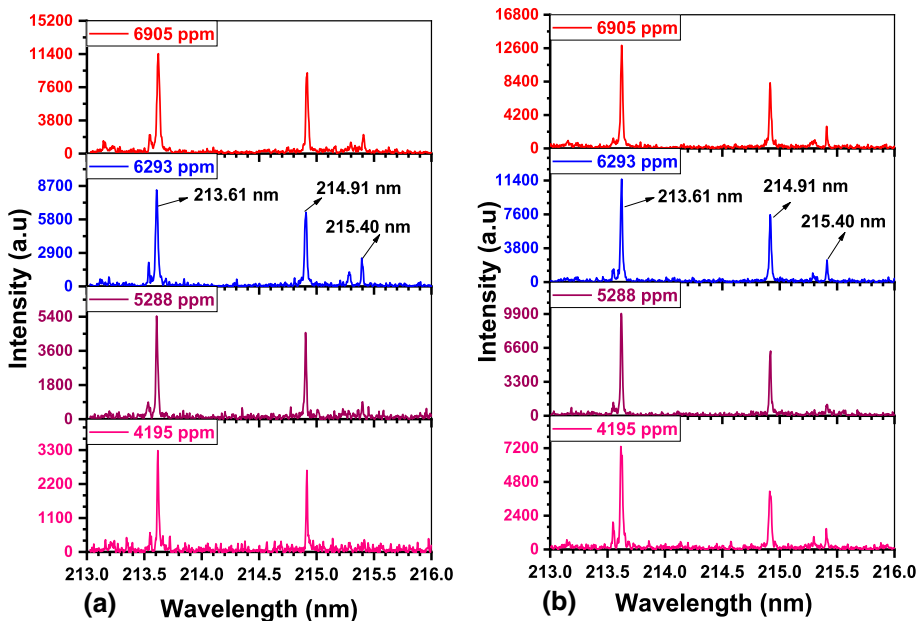
After performing an optimization procedure for maximizing spectral line intensity, the gate width and delay time have been chosen. The choice of the optimum gate width and delay time values for the spectroscopic data acquisition was accomplished by a computer-controlled system. The gate and delay times were adjusted at 1.5 and 0.8  $\mu$ s, respectively. Spectrums from 50 laser shots were averaged to reduce the statistical error due to laser

shot-to-shot fluctuation to improve the LIBS precision. Measurements at new locations on the sample surface were taken to avoid problems linked to sample heterogeneity. Fifty laser shots were fired at each location and saved in separate files, and the average was computed and saved to serve as the library for each spectrum. The plasma emission spectra were collected in both air and helium environments by purging the sample with a helium gas flow. These processes take just a few seconds to identify the whole LIBS spectrum for the studied elemental analysis.

### 3 Result and discussion

#### 3.1 LIBS spectrum studies

The qualitative analysis of the phosphogypsum samples was observed from the optical spectra established from the LIBS technique using 50 laser accumulated shots of NIR nanosecond Nd: YAG pulsed laser with 5 ns pulse duration and 5 mJ pulse energy, detected on 1.5  $\mu$ s gate time, 0.8  $\mu$ s delay. Figure 1a,b represents high-resolution LIBS emission spectra of the non-purged and purged with the He gas on phosphogypsum samples with phosphorus concentrations 4195, 5288, 6293, and 6905 ppm, respectively. All the spectra were recorded under identical experimental conditions. The spectral region from 213 – 216 nm is selected because it gathers the highest phosphorus emission lines intensities proposed to be under optically-thin plasma in LTE as a crucial condition for the plasma calculations parameters. The three selected lines of singly ionized phosphorus (P I) are 213.61, 214.91 and 215.40 nm due to the transitions ( $3s^23p^3, ^2D_{5/2} \rightarrow 3s^23p^2(^3P)4s, ^2P_{3/2}$ ),



**Fig. 1** LIBS emission spectra using NIR laser 1064 nm at pulse energies of 50 mJ of different phosphogypsum samples **a** non-purged and **b** purge

**Table 1** Spectroscopic parameters of the emission lines of P I used to calculate the plasma parameters, taken from the NIST database (Kramida and Ralchenko 2018).

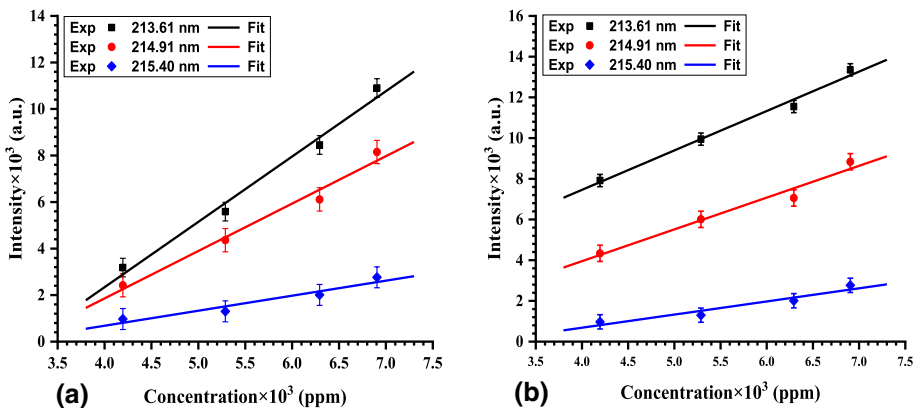
Wavelength (nm) of P I	Transitions		Transition Probability ( $A_{ki}$ ) ( $10^8 s^{-1}$ )	Energies (eV)			$g_k$
	Upper level	Lower level		$E_k$	$E_i$	$\Delta E$	
213.61	$3s^2 3p^3 \ ^2D_{5/2}$	$\longrightarrow$ $3s^2 3p^2(^3P)4s, \ ^2P_{3/2}$	2.83	7.212	1.410	3.816	4
214.91	$3s^2 3p^3 \ ^2D_{3/2}$	$\longrightarrow$ $3s^2 3p^2(^3P)4s, \ ^2P_{1/2}$	3.18	7.175	1.408	3.785	2
215.40	$3s^2 3p^3 \ ^2P_{3/2}$	$\longrightarrow$ $3s^2 3p^2(^1D)4s, \ ^2D_{5/2}$	0.58	8.078	3.816	2.376	6

( $3s^2 3p^3 \ ^2D_{3/2} \rightarrow 3s^2 3p^2(^3P)4s, \ ^2P_{1/2}$ ), and ( $3s^2 3p^3 \ ^2P_{3/2} \rightarrow 3s^2 3p^2(^1D)4s, \ ^2D_{5/2}$ ), respectively. These lines were selected to study the spectroscopic breakdown parameters since they showed the highest correlation values covering the UV wavelength range 213 to 216 nm (Marangoni et al. 2016).

The transitions' parameters of the selected P emission lines have been recognized with the NIST database's help, as demonstrated in Table 1. These data were used to evaluate the spectroscopic breakdown parameters (Kramida and Ralchenko 2018).

### 3.2 Concentration calibration studies

The LIBS concentration calibration curves of the P I spectral lines 213.61, 214.91, and 215.40 nm for phosphogypsum samples in air and purged with He gas at NIR (1064 nm) laser pulse energy 50 mJ are shown in Fig. 2. The calibration curves validate the linear increase in the LIBS signal intensity of the different LIBS spectra wavelengths with the increase of the P concentration at a high correlation coefficient  $\sim 0.975 \pm 0.05$ . The increase in the concentration of the P in the target ablated volume leads to an increase in the ablated mass (Galbács 2017; Qi and Lai 2012; Raciukaitis et al. 2008), which produced an increase in the absorbed laser pulse energy as given by Eq. (1), which is a semi-empirical model given by of G. Raciukaitis et al., in addition to H. Qi and H. Lai (Qi and Lai 2012; Raciukaitis et al. 2008).



**Fig. 2** LIBS concentration calibration curve of P in different phosphogypsum samples **a** Non-purged and **b** Purge using NIR laser 1064 nm at pulse energies of 50 mJ

$$m = \rho V = \rho \frac{\pi w^2 d_0}{4} \left( \ln \frac{2E_p}{\pi w^2 F_{p,c}} \right)^2 \quad (1)$$

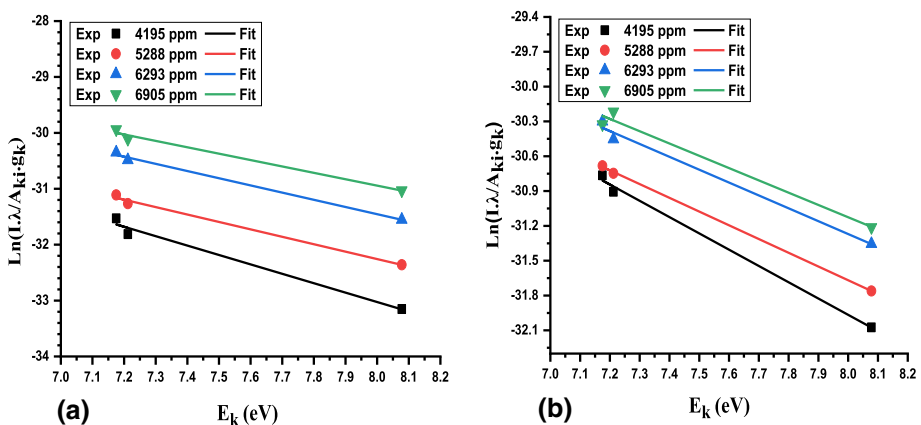
where  $V$  is the ablation volume per a single shot pulse,  $\rho$  is the solid target density,  $w$  is the beam waist of the laser beam at the target surface,  $d_0$  is the effective absorption depth of the target,  $E_p$  is the laser pulse energy,  $F_{p,c}$  is the critical threshold fluence of the laser.

It worth noting that it has been found that in LIBS experiments under LTE conditions, an increase in the target ablated mass may lead to an increase in the plasma electron temperature and density (Fikry et al. 2020b, 2020a). This plasma profile will be verified for LIBS of phosphogypsum samples under the studied conditions.

As shown in Fig. 2, the observed linear calibration relation confirms the absence of the plasma self-absorption and assures the plasma's homogeneity (Palleschi 2020; Sabsabi 2007). The purged samples with the He gas indicated the same trend with an increase of the P concentration at the same laser pulse energy with an enhancement of the signal intensity, especially at the law sample concentrations (Kaski et al. 2003). The small concentrations enhanced approximately to the dappled (Li et al. 2018).

### 3.3 Breakdown spectroscopy parameters and calibration-free studies

The plasma electron temperature and density are considered as the most critical parameters of the breakdown spectroscopy studies (Ahmed et al. 2020; Farooq et al. 2020; Farooq et al. 2018; Farooq et al. 2015; Tawfik et al. 2015b). The plasma temperature is measured under the local thermodynamic equilibrium and optically-thin plasma conditions depending on the previously obtained results of the concentration calibration curve. These conditions will be confirmed later. According to these conditions, the intensities of different fertilisers' emission lines, considered from electron collisions, are much larger than from the radiative processes (Gojani 2012). It is worth noting that the reabsorption effects of plasma emission are negligible (Liu et al. 1999; El Sherbini et al. 2005). Under the LTE conditions, the Boltzmann plot method is given by (Alhijry et al. 2020; Liu et al. 1999);



**Fig. 3** A Boltzmann plots for the P I lines (213.61, 214.91, 215.40 nm) in different concentrations of phosphogypsum samples **a** Non-purged and **b** Purged using NIR laser 1064 nm at pulse energies of 50 mJ

$$\ln \frac{I\lambda}{A_{ki}g_k} = -\frac{1}{KT_e}E_k + \frac{FC}{U(T)} \quad (2)$$

where  $I$  is the intensity of the spectral line,  $\lambda$  is the wavelength of a spectral line,  $K$  is the Boltzmann constant,  $U(T)$  is the partition function,  $A_{ki}$  is the transition probability,  $g_k$  is the statistical weight for the upper level,  $E_k$  is the excited level energy,  $T_e$  is the temperature,  $F$  is an experimental factor and  $C$  is the species concentration.

Figure 3 represents Boltzmann plots for non-purged and purged samples by helium gas for P I lines (213.61, 214.91, and 215.40 nm) at 1064 nm laser wavelength and pulse energy 50 mJ. The  $\ln(I\lambda/A_{ki}g_k)$  is considered for each exciting level energy  $E_k$  under a gradual increase in the P concentration values from 19,200 to 3800 ppm. From Eq. (2), the electron temperature can be determined from each line's slope in Fig. 3, with an uncertainty of about  $\pm 5\%$  (Fikry et al. 2020b, 2020a).

The plasma electron density is considered for the Stark-broadening profile of P I line 213.61 nm because it has the highest correlation coefficient ( $\sim 0.98$ ) with the concentration. The linewidth  $\Delta\lambda_{FWHM}$  is calculated by deconvolution of the phosphorus spectral line profile as a Voigt profile (Alhijry et al. 2020; Fikry et al. 2020a) using Origin software version 9.5 at fixed laser pulse energy at 50 mJ as represented in Fig. 4.

As indicated in Fig. 4, the spectral line FWHM of the P 213.61 nm ranged from 0.02209 to 0.03588 nm with an average error  $\pm 0.00175$  nm and from 0.02337 to 0.03653 with an average  $\pm 0.00125$  nm for the non-purged and purge, respectively. The spectral line FWHM of the P I 213.61 nm under purged with He gas is larger than the non-purged condition, indicating that the purge enhanced the electron density, which verifies by calculating the plasma electron density.

The electron density is measured by the Boltzmann distribution of the electron density as in following in Eq. (3), considering the LTE conditions as referred before (Fikry et al. 2020b; Mortazavi et al. 2014):

$$N_e \approx \left( \frac{\Delta\lambda_{FWHM}}{2W_e} \right) \cdot 10^{16} \quad (3)$$

where  $N_e$  the electron density (in  $\text{cm}^{-3}$ ),  $\Delta\lambda_{FWHM}$  the fundamental line width at half maximum, and  $W_e$  the electron Stark-broadening value. The average value of  $W_e$  for P I 213.61 nm was obtained from H. R. Griem's book (1964) as 0.00295 Å (Griem 1964). The Stark line width  $\Delta\lambda_{FWHM}$  can be corrected by subtracting the instrumental  $\Delta\lambda_{instrument}$  ( $\sim 15$  pm) from the observed line width  $\Delta\lambda_{observed}$  as follows:

$$\Delta\lambda_{FWHM} = \Delta\lambda_{observed} - \Delta\lambda_{instrument} \quad (4)$$

The electron temperature and the electron density calibration curves are directly proportional to the P concentration for the phosphogypsum samples under non-purged and purged He conditions, as shown in Fig. 5. The electron temperatures are found to be ranged from about 6900 to 10,000 K and from 8200 to 11,000 K, for the non-purged, and purged samples, respectively. On the other hand, the electron density of the non-purged and purged samples were found to range from about  $1.1 \times 10^{17}$  to  $3.4 \times 10^{17} \text{ cm}^{-3}$  and from  $1.4 \times 10^{17}$  to  $3.5 \times 10^{17} \text{ cm}^{-3}$ , respectively.

These observations indicated an increase in the plasma electron temperature and density with the phosphorus concentration in the target samples due to the ablation mass increment, which agrees with the previous foundations (Fikry et al. 2020b, 2020a; Galbács 2017; Qi and Lai 2012; Raciukaitis et al. 2008), as discussed above in Eq. (1). The above



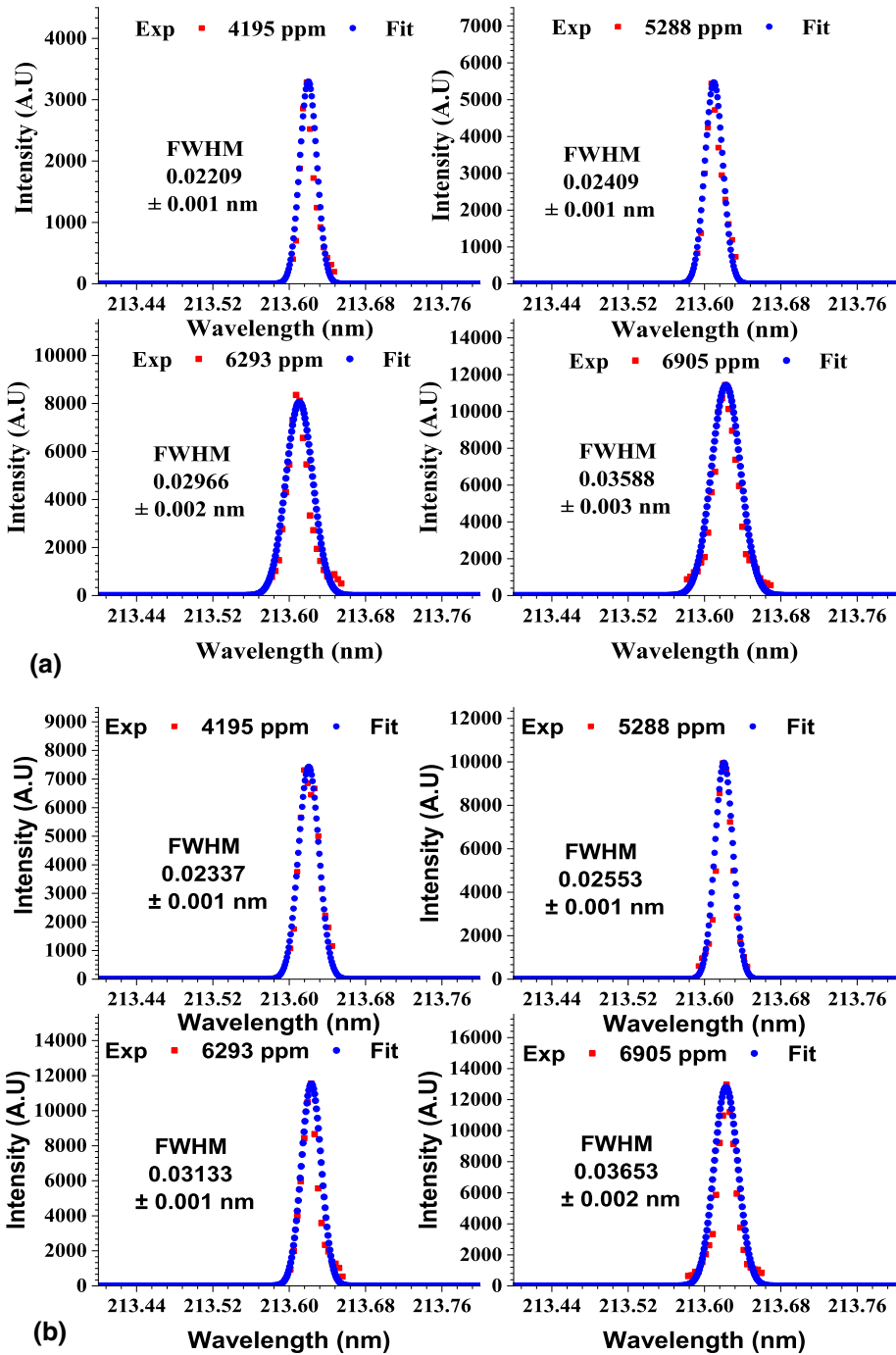
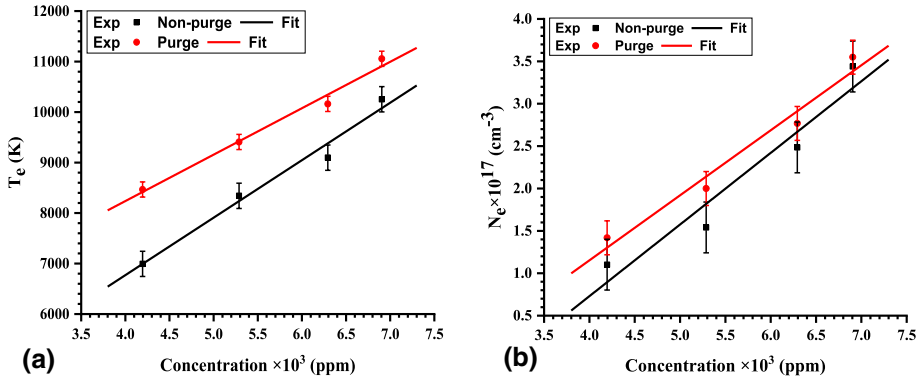


Fig. 4 Voigt Line profile of P I line 213.61 nm in different concentration phosphogypsum samples **a** Non-purged and **b** Purged using NIR laser 1064 nm at pulse energies of 50 mJ



**Fig. 5** LIBS plasma electron temperature ( $T_e$ ) **a** and electron density ( $N_e$ ) at 213.61 nm **b** calibration curves for the P I lines in different phosphogypsum samples non-purged and purged using NIR laser 1064 nm at pulse energies of 50 mJ

results clarify that the ambient plasma conditions (Air or He) play an essential role in the plasma emission characteristics. The plasma expansion rate and the background absorption for the laser energy depend on the ionization energy of the plasma surrounded gas. The main air gases are oxygen and nitrogen, which have ionization energies 13.62 and 14.53 eV, respectively, which are relatively smaller compared with 24.58 eV for He gas (Camacho et al. 2011; Eseller et al. 2012; Kim and Desclaux 2002; Unnikrishnan et al. 2010; Welander 1898; Zavilopulo et al. 2005). Accordingly, then in the case of non-purge (air), the background absorbed more laser energy than in the case of purge (He) (Akram et al. 2017; Effenberger and Scott 2010; Farooq et al. 2014; Rajavelu et al. 2020). These results facilitate the opportunity to directly identify the phosphorus concentration in any unknown phosphogypsum sample by determining  $T_e$  and  $N_e$  values for that sample; then, the P concentration can be observed directly from linear curves in Fig. 5 without needing to complete analysis of the sample or building its calibration curves. The latter represents a novel calibration-free LIBS method to identify the P concentration in unknown phosphogypsum waste samples.

Finally, the verify of the LTE condition, the minimum critical electron density, proposed by McWhirter Eq. (5), which consider the collisional processes are dominated over the radiative processes, has been considered as following (Fikry et al. 2020a; Liu et al. 1999):

$$N_e \geq 1.6 \times 10^{12} \times \Delta E^3 \times T_e^{(1/2)} \tag{5}$$

where  $\Delta E$  is the highest energy difference between the upper and the lower energy level (from Table 1  $\Delta E=3.816$  eV) and  $T_e$  is the plasma temperature (from Fig. 5 for non-purged  $T_e=10,153.1$  K, purged  $T_e=10,955.3$  K). The RHS of Eq. (5), represents the critical  $N_e$   $8.9 \times 10^{15}$   $\text{cm}^{-3}$  for non-purged and  $9.3 \times 10^{15}$   $\text{cm}^{-3}$  for the purged samples. On the other hand, the P I line's observed electron density values were in the range of  $10^{17}$   $\text{cm}^{-3}$  as shown in Fig. 5. Thus, the McWhirter condition is valid, which indicates that the studied plasma can be considered in the LTE condition.

From the linear fitting results of Fig. 5, new empirical formulas can be derived with a high correlation coefficient  $\sim (0.975 - 0.985) \pm 0.05$ . These formulas can be used to identify the unknown P concentration by estimating the plasma parameters ( $T_e$  and  $N_e$ ) using the LIBS method under the experiment mentioned above conditions (1064 nm Nd: YAG laser,

50 mJ pulse energy, 5 ns pulse width, 30 Hz repetition rate, 1.5  $\mu\text{s}$  gate time, 0.8  $\mu\text{s}$  delay time, 50 laser shots, and at the LTE condition). Equations 6 and 7 for the non-purged, 8 and 9 for the He purged which given as follows;

for the non-purged P samples

$$T_e(K) \simeq 2.22 \times 10^3 + 1.14 \times 10^3 \times C(ppm) \quad (6)$$

$$N_e(\text{cm}^{-3}) \times 10^{17} \simeq -2.65 + 0.845 \times C(ppm) \quad (7)$$

For the He purged samples

$$T_e(K) \simeq 4.56 \times 10^3 + 9.19 \times 10^2 \times C(ppm) \quad (8)$$

$$N_e(\text{cm}^{-3}) \times 10^{17} \simeq -1.92 + 0.767 \times C(ppm) \quad (9)$$

## 4 Conclusion

The calibration-free LIBS determined P concentration in any unknown phosphogypsum waste sample directly in the air and helium environment. The plasma electron temperature and density for the known phosphorus concentrations in different phosphogypsum samples were studied. Helium gas purging of the target surface enhanced the LIBS spectrum, plasma electron temperature, and electron density. The increase in the phosphorus concentration in the phosphogypsum samples from 4195 to 6905 ppm change electron temperature values ranged from about 6900 to 10,000 K and from 8200 to 11,000 K for samples with non-purged and purged with helium conditions, respectively. While the electron density values ranged from  $1.1 \times 10^{17}$  to  $3.4 \times 10^{17} \text{ cm}^{-3}$  and  $1.4 \times 10^{17}$  to  $3.5 \times 10^{17} \text{ cm}^{-3}$  for samples under non-purged (air) and He purged conditions, respectively. The observed results indicate that it is possible to improve the exploitation of calibration-free LIBS in the online environmental monitoring by following up only plasma parameters of the phosphogypsum to identify phosphorus concentrations without needing to thoroughly analyze the phosphogypsum, which saves a lot of time and effort.

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